

A lattice spanning-tree entropy function

This article has been downloaded from IOPscience. Please scroll down to see the full text article.

2005 J. Phys. A: Math. Gen. 38 L471

(<http://iopscience.iop.org/0305-4470/38/25/L02>)

View [the table of contents for this issue](#), or go to the [journal homepage](#) for more

Download details:

IP Address: 171.66.16.92

The article was downloaded on 03/06/2010 at 03:48

Please note that [terms and conditions apply](#).

LETTER TO THE EDITOR

A lattice spanning-tree entropy function**M L Glasser^{1,2} and George Lamb³**¹ Departamento de Física Teórica, Atómica y Óptica, Universidad de Valladolid, 47071 Valladolid, Spain² Center for Quantum Device Technology, Clarkson University, Potsdam, NY 13699-5820, USA³ 2942 Ave. del Conquistador, Tucson, AZ 85749-9304, USA

Received 7 April 2005, in final form 13 May 2005

Published 8 June 2005

Online at stacks.iop.org/JPhysA/38/L471**Abstract**

The function

$$W(a, b) = \int_0^{2\pi} dx \int_0^{2\pi} dy \ln[1 - a \cos x - b \cos y - (1 - a - b) \cos(x + y)]$$

which expresses the spanning-tree entropy for various two-dimensional lattices, for example, is evaluated directly in terms of standard functions. It is applied to derive several limiting values of the triangular lattice Green function.

PACS numbers: 02.30.-f, 05.50.+q

Introduction

The function

$$W(a, b) = \int_0^{2\pi} dx \int_0^{2\pi} dy \ln[1 - a \cos x - b \cos y - c \cos(x + y)]$$

with $a + b + c = 1$ arises frequently in the statistical physics and combinatorics of two-dimensional lattice systems. For example:

(i) $a = b = 1/2$

$$S_{sq} = \frac{\ln(2)}{2\pi^2} + \frac{W(a, b)}{4\pi^2}$$

is the spanning-tree entropy for the square lattice [1].

(ii) $a = b = 1/3$

$$S_{tr} = \frac{\ln(6)}{4\pi^2} + \frac{W(a, b)}{4\pi^2}$$

is the spanning-tree entropy for the triangular lattice [2].

(iii)

$$a = \frac{\sinh K_1}{\sinh K_1 + \sinh K_2 + \sinh K_3}$$

$$b = \frac{\sinh K_2}{\sinh K_1 + \sinh K_2 + \sinh K_3}$$

$$F_I = \ln(2) + \frac{1}{8\pi^2} \ln[\sinh K_1 + \sinh K_2 + \sinh K_3] + \frac{W(a, b)}{8\pi^2}$$

is the critical free energy of the Ising model on a triangular lattice [3].

By comparing the free energies of various Potts models which are known to be related and have been worked out in different ways, Chen and Wu [4] have proposed that

$$\frac{1}{4\pi^2} \int_0^{2\pi} d\theta \int_0^{2\pi} d\phi \ln[A + B + C - A \cos \theta - B \cos \phi - C \cos(\theta + \phi)]$$

$$= -\ln(2S) + \frac{2}{\pi} [Ti_2(AS) + Ti_2(BS) + Ti_2(CS)]$$

where $A, B, C \geq 0$ and $S = 1/\sqrt{AB + BC + CA}$. The aim of this letter is to provide a direct proof of this formula.

Calculation

By symmetry, one easily finds

$$W(a, b) = 2[W_+(a, b) + W_-(a, b)] \quad (1)$$

where

$$W_{\pm}(a, b) = \int_0^{\pi} dx \int_0^{\pi} dy \ln[1 - a \cos x - b \cos y - c \cos(x \pm y)].$$

Next, we make the standard change of variable

$$u = \tan(x/2) \quad v = \tan(y/2)$$

to obtain

$$W_+(a, b) = 4 \int_0^{\infty} \int_0^{\infty} \frac{du dv}{(1+u^2)(1+v^2)} \ln \left[\frac{2}{(1+u^2)(1+v^2)} \right] + 4 \int_0^{\infty} \int_0^{\infty} \ln[(u+v)^2 - a((u+v)^2 - u^2(1+v^2)) - b((u+v)^2 - v^2(1+u^2))] \frac{du dv}{(1+u^2)(1+v^2)}.$$

The first integral is elementary, giving

$$W_+(a, b) = -3\pi^2 \ln(2) + 4 \int_0^{\infty} \int_0^{\infty} \frac{\ln[c(u+v)^2 + au^2(1+v^2) + bv^2(1+u^2)]}{(1+u^2)(1+v^2)} du dv.$$

Similarly,

$$W_-(a, b) = -3\pi^2 \ln(2) + 4 \int_0^{\infty} \int_0^{\infty} \frac{\ln[c(u-v)^2 + au^2(1+v^2) + bv^2(1+u^2)]}{(1+u^2)(1+v^2)} du dv.$$

By inserting the last two expressions into (1) and noting that the resulting integrand is even in v , we have

$$W(a, b) = -12\pi^2 \ln(2) + 8F(a, b),$$

where

$$F(a, b) = \int_0^{\infty} \frac{du}{1+u^2} \int_{-\infty}^{\infty} \frac{dv}{1+v^2} \ln[(1-b)u^2 + (1-a)v^2 + (a+b)u^2v^2 + 2cuv].$$

Next, let $u = vw$ to obtain

$$F(a, b) = 2 \int_0^\infty \frac{u \, du}{1+u^2} \ln(u) \int_{-\infty}^\infty \frac{dw}{1+u^2w^2} + \int_0^\infty \frac{du}{u(1+u^2)} \int_{-\infty}^\infty \frac{dw}{w^2+u^{-2}} \times \ln[1-b+(1-a)w^2+(a+b)u^2w^2+2cw].$$

The first integral vanishes and since [5]

$$\int_0^\infty \frac{dw}{w^2+d^2} \ln[\alpha w^2+2\beta w+\gamma] = \frac{\pi}{d} \ln[\alpha d^2+\gamma+2d\sqrt{\alpha\gamma-\beta^2}],$$

after the substitution $u = 1/z$,

$$F(a, b) = \pi \int_0^\infty \frac{dz}{1+z^2} \ln[(1-a)z^2+(1+a)+2\sqrt{[ac+b(1-b)]z^2+(1-b)(a+b)}]. \tag{2}$$

Now, let us define

$$A = \frac{y \cot(\theta/2)}{1+\sqrt{1+y^2}}, \quad B = \frac{y \tan(\theta/2)}{1+\sqrt{1+y^2}}.$$

Then $(A+B)/(1-AB) = y \csc \theta$, so $\tan^{-1}(y \csc \theta) = \tan^{-1} A + \tan^{-1} B$. However,

$$\csc \theta \tan^{-1}(y \csc \theta) = -\frac{d}{d\theta} \int_1^{\csc \theta} \frac{\tan^{-1}(yu)}{\sqrt{u^2-1}} du$$

$$\csc \theta \tan^{-1} A = -\frac{d}{d\theta} \int_0^A \frac{\tan^{-1} u}{u} du$$

$$\csc \theta \tan^{-1} B = \frac{d}{d\theta} \int_0^B \frac{\tan^{-1} u}{u} du.$$

Hence, since both sides vanish for $\theta = \pi/2$,

$$\int_1^{\csc \theta} \frac{\tan^{-1}(yu)}{\sqrt{u^2-1}} du = Ti_2(A) - Ti_2(B) \tag{3}$$

where

$$Ti_2(z) = \int_0^z \frac{\tan^{-1} x}{x} dx.$$

With

$$\theta = \csc^{-1} \sqrt{\frac{b^2-a^2+1}{1-a^2}}, \quad y = \sqrt{a^{-2}-1}, \quad u = \sqrt{\frac{x^2-a^2+1}{1-a^2}}$$

in (3), one obtains

$$\int_0^b \frac{dx}{\sqrt{x^2-a^2+1}} \tan^{-1} \frac{\sqrt{x^2-a^2+1}}{a} = Ti_2\left(\frac{\sqrt{b^2+1-a^2}+b}{1+a}\right) - Ti_2\left(\frac{\sqrt{b^2+1-a^2}-b}{1+a}\right). \tag{4}$$

We next consider the integral

$$g(a, b) = \int_0^\infty \frac{ds}{s^2+1} \ln[\sqrt{b^2(s^2+1)+1+a}]$$

for which it is elementary to determine

$$g(a, 0) = \frac{\pi}{2} \ln(1+a), \quad \frac{\partial}{\partial u} g(a, u) = \frac{\tan^{-1}(\sqrt{u^2+1-a^2}/a)}{\sqrt{u^2+1-a^2}}.$$

Therefore, by integrating over u using (4), we find

$$g(a, b) = \frac{\pi}{2} \ln(1+a) + Ti_2 \left(\frac{\sqrt{b^2+1-a^2}+b}{1+a} \right) - Ti_2 \left(\frac{\sqrt{b^2+1-a^2}-b}{1+a} \right)$$

which is easily transformed into

$$\int_0^\infty \frac{\ln[\alpha + \sqrt{\beta^2 x^2 + \gamma^2}]}{x^2 + 1} dx = \frac{\pi}{2} \ln[\beta + \sqrt{\gamma^2 - \alpha^2}] + Ti_2 \left(\frac{\alpha + \sqrt{\gamma^2 - \beta^2}}{\beta + \sqrt{\gamma^2 - \alpha^2}} \right) + Ti_2 \left(\frac{\alpha - \sqrt{\gamma^2 - \beta^2}}{\beta + \sqrt{\gamma^2 - \alpha^2}} \right). \quad (5)$$

The argument of the logarithm in the integrand of $F(a, b)$ in (2) can be factored: let $R = \sqrt{[ac + b(1-b)]x^2 + (1-b)(a+b)}$; then

$$(1-a)x^2 + 1 + a + 2R = [ac + b(1-b)]^{-1} [a(1-a) + 2bc + (1-a)R][a + R].$$

The integrals resulting from inserting this into (2) are either elementary or can be evaluated by using (5). After some algebraic manipulation, we obtain

$$W(a, b) = 4\pi^2 \ln \frac{d}{2} + 8\pi [Ti_2(a/d) + Ti_2(b/d) + Ti_2(c/d)] \quad (6)$$

where $d = \sqrt{ab + bc + ac}$, which is equivalent to Chen and Wu's formula.

In conclusion, we list a few values of the anisotropic triangular lattice Green function that can be obtained from (6) by differentiation. Here, $\Delta(a, b, c) = a + b + c - a \cos x - b \cos y - c \cos(x + y)$, $d = \sqrt{ab + bc + ca}$.

$$\begin{aligned} & \int_0^{2\pi} dx \int_0^{2\pi} dy \frac{\Delta(1, 0, 0)}{\Delta(a, b, c)} \\ &= \frac{4\pi}{d^2} (b+c) \left[\frac{\pi}{2} + \frac{a(b+c) + 2bc}{a(b+c)} \tan^{-1}(a/d) - \tan^{-1} \left(\frac{(b+c)d}{bc+d^2} \right) \right] \\ & \int_0^{2\pi} dx \int_0^{2\pi} dy \frac{\Delta(0, 0, 1)}{\Delta(a, b, c)} \\ &= \frac{4\pi}{d^2} (a+b) \left[\frac{\pi}{2} + \frac{(a+b)c + 2ab}{(a+b)c} \tan^{-1}(c/d) - \tan^{-1} \left(\frac{(a+b)d}{ab+d^2} \right) \right] \\ & \int_0^{2\pi} dx \int_0^{2\pi} dy \frac{\Delta(1, 1, 1)}{\Delta(a, b, c)} = \frac{4\pi^2}{d^2} (a+b+c) \\ & \quad + \frac{8\pi}{d^2} \left[\frac{bc-a^2}{a} \tan^{-1}(a/d) + \frac{ac-b^2}{b} \tan^{-1}(b/d) + \frac{ab-c^2}{c} \tan^{-1}(c/d) \right] \\ & \int_0^{2\pi} dx \int_0^{2\pi} dy \frac{\Delta(-1, 1, 0)}{\Delta(a, b, c)} = \frac{2\pi^2}{d^2} (a-b) + \frac{4\pi}{d^2} \left[(b-a) \tan^{-1}(c/d) \right. \\ & \quad \left. + \left(a+b + \frac{2c(a+b)}{b} \right) \tan^{-1}(b/d) - \left(a+b+2c + \frac{2c(a+b)}{a} \right) \tan^{-1}(a/d) \right]. \end{aligned}$$

Acknowledgments

This work was supported in part by the Spanish MEC (BFM2002-03773 and MLG grant SAB2003-0117) and Junta de Castilla y Leon (VA085/02). MLG thanks the Universidad de Valladolid for hospitality and the NSF (USA) for partial support (DMR-0121146).

References

- [1] Wu F Y 1977 *J. Phys. A: Math. Gen.* **10** L113
- [2] Glasser M L and Wu F Y 2003 *Preprint* cond-matter/0309198 (*Ramanujan J.* to appear)
- [3] Houtappel R M F 1950 *Physica* **16** 425
- [4] Chen L C and Wu F Y 2005 *Preprint* cond-matter/0501228
- [5] Lewin L 1958 *Dilogarithms and Associated Functions* (London: MacDonald) p 268